## Krein Signatures for Eigenvalue Problems Associated to Integrable Systems

Mathew Johnson - UIUC Mathematics
The XIXth International Workshop on Operator Theory and its Applications
College of William \& Mary

$$
7 / 24 / 2008
$$

(Joint work with Jared Bronski - UIUC Mathematics)

## Abstract

There is a result of Klaus and Shaw which shows that the Zakharov-Shabat eigenvalue problem has discrete spectrum which lies on the imaginary axis if the potential has a single critical point and decays monotonically away from this point (the potential is monomodal). We put this calculation in the context of the Krein signature (a tool for studying the stability of symplectic matrices) and prove an analogous theorem for the eigenvalue problem which solves the Sine-Gordon equation (in laboratory coordinates). This is joint work with Jared Bronski.

## Outline

(1) Introduction

- Preliminaries
- Known Results
- Motivation
(2) Krein Signatures
- Symmetries \& Setup
- Application to SG Eigenvalue Problem
(3) Confinement Results
- Topological Charge $\pm 1$
- Topological Charge 0

4) Extensions

## Introduction

- We consider the Cauchy problem for the Sine-Gordon equation in laboratory coordinates:

$$
\begin{aligned}
u_{x x}-u_{t t} & =\sin u \\
u(x, 0) & =u(x) \\
u_{t}(x, 0) & =v(x)
\end{aligned}
$$

## Introduction

- We consider the Cauchy problem for the Sine-Gordon equation in laboratory coordinates:

$$
\begin{aligned}
u_{x x}-u_{t t} & =\sin u \\
u(x, 0) & =u(x) \\
u_{t}(x, 0) & =v(x)
\end{aligned}
$$

- Toy Model: Consider an infinite wire with a continuum of coupled pendulums, with the pendulum at position $x$ at time $t$ hanging with angle $u(x, t)$ with respect to its rest position.


## Introduction

- Assumptions on $u: u \in C^{1}(\mathbb{R}), \lim _{x \rightarrow \pm \infty} u(x)=2 \pi k_{ \pm}$


## Introduction

- Assumptions on $u: u \in C^{1}(\mathbb{R}), \lim _{x \rightarrow \pm \infty} u(x)=2 \pi k_{ \pm}$
- Define $Q_{\text {top }}:=\frac{1}{2 \pi} \int_{-\infty}^{\infty} u^{\prime}(x) d x=k_{+}-k_{-}$to be the "topological charge" of the potential $u$.


## Introduction

- Assumptions on $u: u \in C^{1}(\mathbb{R}), \lim _{x \rightarrow \pm \infty} u(x)=2 \pi k_{ \pm}$
- Define $Q_{\text {top }}:=\frac{1}{2 \pi} \int_{-\infty}^{\infty} u^{\prime}(x) d x=k_{+}-k_{-}$to be the "topological charge" of the potential $u$.
- We consider only initial data with topological charge 0 (breathers) or $\pm 1$ (kinks).


## Introduction

- This equation is integrable via the Inverse Scattering Transform with corresponding spectral problem (Faddeev-Tahktajan / Kaup)

$$
\begin{equation*}
4 \Phi_{x}=\left(z-\frac{1}{z}\right) \tau_{1} \cos \left(\frac{u}{2}\right) \Phi+\left(z+\frac{1}{z}\right) \tau_{2} \sin \left(\frac{u}{2}\right) \Phi-v \tau_{3} \Phi \tag{1}
\end{equation*}
$$

considered on $L^{2}(d x ; \mathbb{C})$, where $z \in \mathbb{C}, \Phi=\binom{\phi_{1}}{\phi_{2}}$,

- $u=u(x, 0)$
- $v=u_{t}(x, 0)$
and

$$
\tau_{1}=\left(\begin{array}{cc}
-i & 0 \\
0 & i
\end{array}\right), \quad \tau_{2}=\left(\begin{array}{cc}
0 & i \\
i & 0
\end{array}\right), \quad \tau_{3}=\left(\begin{array}{cc}
0 & 1 \\
-1 & 0
\end{array}\right)
$$

## Introduction

- This equation is integrable via the Inverse Scattering Transform with corresponding spectral problem (Faddeev-Tahktajan / Kaup)

$$
\begin{equation*}
4 \Phi_{x}=\left(z-\frac{1}{z}\right) \tau_{1} \cos \left(\frac{u}{2}\right) \Phi+\left(z+\frac{1}{z}\right) \tau_{2} \sin \left(\frac{u}{2}\right) \Phi-v \tau_{3} \Phi \tag{1}
\end{equation*}
$$

considered on $L^{2}(d x ; \mathbb{C})$, where $z \in \mathbb{C}, \Phi=\binom{\phi_{1}}{\phi_{2}}$,

- $u=u(x, 0)$
- $v=u_{t}(x, 0)$
and

$$
\tau_{1}=\left(\begin{array}{cc}
-i & 0 \\
0 & i
\end{array}\right), \quad \tau_{2}=\left(\begin{array}{cc}
0 & i \\
i & 0
\end{array}\right), \quad \tau_{3}=\left(\begin{array}{cc}
0 & 1 \\
-1 & 0
\end{array}\right)
$$

- "Main Result": Under certain conditions on $u$ and $v$, one must have $z \in S^{1}$.


## Known Results

- Characteristic coordinates $\chi=\frac{x+t}{\sqrt{2}}, \eta=\frac{x-t}{\sqrt{2}}$

$$
\Rightarrow u_{\chi \eta}=\sin (u)
$$

This is the isospectral flow for the Zhakarov-Shabat system

$$
\begin{aligned}
& v_{1, \chi}=-i z v_{1}+q v_{2} \\
& v_{2, \chi}=i z v_{2}-q^{*} v_{1}
\end{aligned}
$$

where $q$ is related to the initial data, $z \in \mathbb{C}$, and $*$ denotes complex conjugation.

## Known Results

- Characteristic coordinates $\chi=\frac{x+t}{\sqrt{2}}, \eta=\frac{x-t}{\sqrt{2}}$

$$
\Rightarrow u_{\chi \eta}=\sin (u)
$$

This is the isospectral flow for the Zhakarov-Shabat system

$$
\begin{aligned}
v_{1, \chi} & =-i z v_{1}+q v_{2} \\
v_{2, \chi} & =i z v_{2}-q^{*} v_{1}
\end{aligned}
$$

where $q$ is related to the initial data, $z \in \mathbb{C}$, and $*$ denotes complex conjugation.

- We have


## Theorem (Klaus \& Shaw $(2001,2002)$ )

If $q: \mathbb{R} \rightarrow \mathbb{R}$ is in $L^{1} \cap C^{1}$ and has only one critical point (a local max), then the discrete spectrum lies on the imaginary axis in the $z$-plane.

## Known Results

- FACT: that the phase of the potential $q$ is related to the momentum of the initial pulse, with zero phase corresponding to with zero initial momentum, i.e. stationary initial data.


## Known Results

- FACT: that the phase of the potential $q$ is related to the momentum of the initial pulse, with zero phase corresponding to with zero initial momentum, i.e. stationary initial data.
- FACT: For Z-S system, eigenvalues on $\mathbb{R} i$ correspond to stationary solitons.


## Known Results

- FACT: that the phase of the potential $q$ is related to the momentum of the initial pulse, with zero phase corresponding to with zero initial momentum, i.e. stationary initial data.
- FACT: For Z-S system, eigenvalues on $\mathbb{R} i$ correspond to stationary solitons.
- Klaus-Shaw result states that, under certain monotonicity conditions, a pulse with zero initial momentum give rise to solitons with zero momentum.


## Known Results

- FACT: that the phase of the potential $q$ is related to the momentum of the initial pulse, with zero phase corresponding to with zero initial momentum, i.e. stationary initial data.
- FACT: For Z-S system, eigenvalues on $\mathbb{R} i$ correspond to stationary solitons.
- Klaus-Shaw result states that, under certain monotonicity conditions, a pulse with zero initial momentum give rise to solitons with zero momentum.
- In laboratory coordinates, initial data with zero initial momentum corresponds to $v(x)=u_{t}(x, 0)=0$ and stationary solitons correspond to $z \in S^{1}$.


## Known Results

- FACT: that the phase of the potential $q$ is related to the momentum of the initial pulse, with zero phase corresponding to with zero initial momentum, i.e. stationary initial data.
- FACT: For Z-S system, eigenvalues on $\mathbb{R} i$ correspond to stationary solitons.
- Klaus-Shaw result states that, under certain monotonicity conditions, a pulse with zero initial momentum give rise to solitons with zero momentum.
- In laboratory coordinates, initial data with zero initial momentum corresponds to $v(x)=u_{t}(x, 0)=0$ and stationary solitons correspond to $z \in S^{1}$.
- Buckingham \& Miller (2008): If one takes initial data which satisfies

$$
\sin (u(x) / 2)=\operatorname{sech} x, \quad \cos (u(x) / 2)=\tanh (x), \quad v(x)=0
$$

eigenvalues are confined to unit circle and are simple. Notice such a potential is monotone with $Q_{t o p}=-1$.

## Main Theorem

- Recall (1) is given by

$$
4 \Phi_{x}=\left(z-\frac{1}{z}\right) \tau_{1} \cos \left(\frac{u}{2}\right) \Phi+\left(z+\frac{1}{z}\right) \tau_{2} \sin \left(\frac{u}{2}\right)-v \tau_{3} \Phi
$$

## Theorem (Bronski \& M.J.)

Consider the Sine-Gordon equation in laboratory coordinates. Let $v=u_{t}(x, 0)=0$ and let $u=u(x, 0)$ be such that one of the following conditions holds:
(1) $u$ is monotone with $Q_{\text {top }}= \pm 1$, or
(2) $u$ has $Q_{\text {top }}=0$ and is monomodal with a positive maximum $u_{0} \in(0, \pi)$.

Then the discrete spectrum of (1) is simple and lies on the unit circle.

## Main Theorem

- Recall (1) is given by

$$
4 \Phi_{x}=\left(z-\frac{1}{z}\right) \tau_{1} \cos \left(\frac{u}{2}\right) \Phi+\left(z+\frac{1}{z}\right) \tau_{2} \sin \left(\frac{u}{2}\right)-v \tau_{3} \Phi
$$

## Theorem (Bronski \& M.J.)

Consider the Sine-Gordon equation in laboratory coordinates. Let $v=u_{t}(x, 0)=0$ and let $u=u(x, 0)$ be such that one of the following conditions holds:
(1) $u$ is monotone with $Q_{\text {top }}= \pm 1$, or
(2) $u$ has $Q_{\text {top }}=0$ and is monomodal with a positive maximum $u_{0} \in(0, \pi)$.

Then the discrete spectrum of (1) is simple and lies on the unit circle.

- Main analytical tool: Krein signatures, i.e. functionals on $L^{2}$ which encode information about eigenvalues and eigenspaces.


## Symmetries

As we will see, Krein signatures are built in such a way to abuse the symmetries of the discrete spectrum. So, we begin with the following lemma:

## Lemma

The symmetry group of the discrete spectrum of (1) is $\mathbb{Z}_{2} \times \mathbb{Z}_{2} \times \mathbb{Z}_{2}$, corresponding to reflection across the real and imaginary axis, as well as the unit circle.


## Krein Signatures

- Suppose we can write (1) as

$$
M v=z v
$$

Symmetries $\Rightarrow$ " $M \sim M^{-\dagger " . ~}$

## Krein Signatures

- Suppose we can write (1) as

$$
M v=z v
$$

Symmetries $\Rightarrow$ " $M \sim M^{-\dagger " . ~}$

- If $U$ is such that $M^{\dagger} U M=U$, define the associated KREIN SIGNAGURE to be

$$
\kappa(v):=\langle v, U v\rangle
$$

for $v \in L^{2}$.

## Krein Signatures

- Suppose we can write (1) as

$$
M v=z v
$$

Symmetries $\Rightarrow$ " $M \sim M^{-\dagger " . ~}$

- If $U$ is such that $M^{\dagger} U M=U$, define the associated KREIN SIGNAGURE to be

$$
\kappa(v):=\langle v, U v\rangle
$$

for $v \in L^{2}$.

- Basic Result: If $M v=\lambda v$,

$$
\kappa(v) \neq 0 \Rightarrow \lambda \in S^{1}
$$

## Krein Signatures

- PROOF:

$$
\lambda \kappa(v)=\langle v, U \lambda v\rangle=\langle v, U M v\rangle
$$

## Krein Signatures

- PROOF:

$$
\begin{aligned}
\lambda \kappa(v) & =\langle v, U \lambda v\rangle=\langle v, U M v\rangle \\
& =\left\langle v, M^{-\dagger} U v\right\rangle
\end{aligned}
$$

## Krein Signatures

- PROOF:

$$
\begin{aligned}
\lambda \kappa(v) & =\langle v, U \lambda v\rangle=\langle v, U M v\rangle \\
& =\left\langle v, M^{-\dagger} U v\right\rangle \\
& =\left\langle M^{-1} v, U v\right\rangle
\end{aligned}
$$

## Krein Signatures

- PROOF:

$$
\begin{aligned}
\lambda \kappa(v) & =\langle v, U \lambda v\rangle=\langle v, U M v\rangle \\
& =\left\langle v, M^{-\dagger} U v\right\rangle \\
& =\left\langle M^{-1} v, U v\right\rangle \\
& =(\bar{\lambda})^{-1} \kappa(v)
\end{aligned}
$$

## Krein Signatures

- PROOF:

$$
\begin{aligned}
\lambda \kappa(v) & =\langle v, U \lambda v\rangle=\langle v, U M v\rangle \\
& =\left\langle v, M^{-\dagger} U v\right\rangle \\
& =\left\langle M^{-1} v, U v\right\rangle \\
& =(\bar{\lambda})^{-1} \kappa(v)
\end{aligned}
$$

- Thus,

$$
\left(\lambda-\frac{1}{\bar{\lambda}}\right) \kappa(v)=0
$$

which proves the claim.

## Strategy for Confinement

- Determine the symmetries of the discrete spectrum of (1).


## Strategy for Confinement

- Determine the symmetries of the discrete spectrum of (1).
- Construct Krein signature $\kappa$ corresponding to reflection about $S^{1}$.


## Strategy for Confinement

- Determine the symmetries of the discrete spectrum of (1).
- Construct Krein signature $\kappa$ corresponding to reflection about $S^{1}$.
- Prove $\kappa(v) \neq 0$ for any eigenfunction $v$.


## More Krein Signatures

- Moreover, $\kappa(v) \neq 0$ implies eigenspace is semi-simple, i.e. the eigenspace is diagonalizable.


## More Krein Signatures

- Moreover, $\kappa(v) \neq 0$ implies eigenspace is semi-simple, i.e. the eigenspace is diagonalizable.
- Claim: For 2nd order ODE eigenvalue problems, semi-simple implies simple.


## More Krein Signatures

- Moreover, $\kappa(v) \neq 0$ implies eigenspace is semi-simple, i.e. the eigenspace is diagonalizable.
- Claim: For 2nd order ODE eigenvalue problems, semi-simple implies simple.
- Proof: The existence of a semi-simple eigenspace of multiplicity higher than one implies $\exists$ two linearly independent solutions decaying as $x \rightarrow \infty$. This contradicts the asymptotic behavior of the Jost solutions. $\square$


## Example

- It is possible to generalize the above theory to different spectral symmetries.


## Example

- It is possible to generalize the above theory to different spectral symmetries.
- The Zakharov-Shabat eigenvalue problem (for real potentials) is given by $M v=\lambda v$ on $L^{2}(d x ; \mathbb{R})$, where

$$
M=\left(\begin{array}{cc}
i \partial_{x} & -i q(x) \\
-i q(x) & -i \partial_{x}
\end{array}\right) \Rightarrow M \sim-M^{\dagger}
$$

## Example

- It is possible to generalize the above theory to different spectral symmetries.
- The Zakharov-Shabat eigenvalue problem (for real potentials) is given by $M v=\lambda v$ on $L^{2}(d x ; \mathbb{R})$, where

$$
M=\left(\begin{array}{cc}
i \partial_{x} & -i q(x) \\
-i q(x) & -i \partial_{x}
\end{array}\right) \Rightarrow M \sim-M^{\dagger}
$$

- $M$ satisfies $U M=-M^{\dagger} U$ for $U=\left(\begin{array}{ll}0 & 1 \\ 1 & 0\end{array}\right)$ (corresponding to reflection of $\operatorname{spec}(M)$ about $\mathbb{R} i)$.


## Example

- It is possible to generalize the above theory to different spectral symmetries.
- The Zakharov-Shabat eigenvalue problem (for real potentials) is given by $M v=\lambda v$ on $L^{2}(d x ; \mathbb{R})$, where

$$
M=\left(\begin{array}{cc}
i \partial_{x} & -i q(x) \\
-i q(x) & -i \partial_{x}
\end{array}\right) \Rightarrow M \sim-M^{\dagger}
$$

- $M$ satisfies $U M=-M^{\dagger} U$ for $U=\left(\begin{array}{ll}0 & 1 \\ 1 & 0\end{array}\right)$ (corresponding to reflection of $\operatorname{spec}(M)$ about $\mathbb{R} i)$.

$$
\kappa=\langle\phi, U \phi\rangle_{L^{2}}=\int_{\mathbb{R}}\left(\phi_{1}^{*} \phi_{2}+\phi_{2}^{*} \phi_{1}\right) d x
$$

This is exactly the quantity Klaus and Shaw consider in their papers.

## Krein Signatures

## Lemma

The eigenvalue problem (1) can be written as $M v=\lambda v$, where the operator $M$ above satisfies $M^{\dagger} U M=U$ where $U$ is of the form

$$
\left(\begin{array}{cccc}
0 & 0 & \cos \left(\frac{u(x)+\pi}{2}\right) & -\sin \left(\frac{u(x)+\pi}{2}\right) \\
0 & 0 & \sin \left(\frac{u(x)+\pi}{2}\right) & \cos \left(\frac{u(x)+\pi}{2}\right) \\
-\cos \left(\frac{u(x)+\pi}{2}\right) & -\sin \left(\frac{u(x)+\pi}{2}\right) & 0 & 0 \\
\sin \left(\frac{u(x)+\pi}{2}\right) & -\cos \left(\frac{u(x)+\pi}{2}\right) & 0 & 0
\end{array}\right)
$$

$$
\begin{aligned}
z=r e^{i \theta} & \Rightarrow \\
\kappa & =i\langle\Phi, U \Phi\rangle_{L^{2}} \\
& =\sin \theta \int_{\mathbb{R}} \sin \left(\frac{u}{2}\right)|\Phi|^{2} d x-i \cos \theta \int_{\mathbb{R}} \cos \left(\frac{u}{2}\right)\left\langle\Phi, \tau_{3} \Phi\right\rangle d x
\end{aligned}
$$

## Topological Charge $\pm 1$ : Confinement Result

- First Main Result:


## Theorem

Let $u(x)$ be a monotone potential satisfying $u(x) \rightarrow 0$ as $x \rightarrow-\infty$ and $u(x) \rightarrow 2 \pi$ as $x \rightarrow \infty$, i.e. $u$ has topological charge $Q_{\text {top }}=1$. Then the discrete spectrum of (1) lies on the unit circle and is simple.

## Topological Charge $\pm 1$ : Confinement Result

- First Main Result:


## Theorem

Let $u(x)$ be a monotone potential satisfying $u(x) \rightarrow 0$ as $x \rightarrow-\infty$ and $u(x) \rightarrow 2 \pi$ as $x \rightarrow \infty$, i.e. $u$ has topological charge $Q_{\text {top }}=1$. Then the discrete spectrum of (1) lies on the unit circle and is simple.

- Proof: Boundary Conditions $\Rightarrow \phi_{2}$ generically grows exponentially as $x \rightarrow \pm \infty$. Suggests we should express $\kappa$ in terms of $\phi_{2}$ :


## Topological Charge $\pm 1$ : Confinement Result

- First Main Result:


## Theorem

Let $u(x)$ be a monotone potential satisfying $u(x) \rightarrow 0$ as $x \rightarrow-\infty$ and $u(x) \rightarrow 2 \pi$ as $x \rightarrow \infty$, i.e. $u$ has topological charge $Q_{\text {top }}=1$. Then the discrete spectrum of (1) lies on the unit circle and is simple.

- Proof: Boundary Conditions $\Rightarrow \phi_{2}$ generically grows exponentially as $x \rightarrow \pm \infty$. Suggests we should express $\kappa$ in terms of $\phi_{2}$ :
- Using original eigenvalue problem, eventually get

$$
\kappa=2\left(r+\frac{1}{r}\right) \int_{\mathbb{R}} \frac{\sin (\theta)}{\sin (u / 2)}\left|\phi_{2}\right|^{2} d x+2 \int_{\mathbb{R}} \frac{u_{x}}{\sin ^{2}(u / 2)}\left|\phi_{2}\right|^{2} d x>0
$$

This completes the proof.

## Topological Charge $\pm 1$ : Confinement Result

- First Main Result:


## Theorem

Let $u(x)$ be a monotone potential satisfying $u(x) \rightarrow 0$ as $x \rightarrow-\infty$ and $u(x) \rightarrow 2 \pi$ as $x \rightarrow \infty$, i.e. $u$ has topological charge $Q_{\text {top }}=1$. Then the discrete spectrum of (1) lies on the unit circle and is simple.

- Proof: Boundary Conditions $\Rightarrow \phi_{2}$ generically grows exponentially as $x \rightarrow \pm \infty$. Suggests we should express $\kappa$ in terms of $\phi_{2}$ :
- Using original eigenvalue problem, eventually get

$$
\kappa=2\left(r+\frac{1}{r}\right) \int_{\mathbb{R}} \frac{\sin (\theta)}{\sin (u / 2)}\left|\phi_{2}\right|^{2} d x+2 \underbrace{\int_{\mathbb{R}} \frac{u_{x}}{\sin ^{2}(u / 2)}\left|\phi_{2}\right|^{2} d x}_{\lim _{x \rightarrow \pm \infty} \phi_{2}(x)=0}>0
$$

This completes the proof.

## Topological Charge $0: 0<\nu_{0} \leq \frac{\pi}{2}$


$\Rightarrow$ All eigenvalues lie on unit circle.

## Topological Charge $0: \frac{\pi}{2}<\nu_{0}<\pi$



## Topological Charge $0: \frac{\pi}{2}<\omega_{0}<\pi$

- We will count the number of eigenvalues in two different ways:
(1) As you vary $u_{0}$, count how many times an eigenvalue either emerges from or gets absorbed into the continuous spectrum (homotopy argument). Yields upper bound.
(2) Count number of points on $S^{1}$ which correspond to an eigenvalue (winding number argument). Clearly gives lower bound.


## Topological Charge 0: $\frac{\pi}{2}<u_{0}<\pi$

- We will count the number of eigenvalues in two different ways:
(1) As you vary $u_{0}$, count how many times an eigenvalue either emerges from or gets absorbed into the continuous spectrum (homotopy argument). Yields upper bound.
(2) Count number of points on $S^{1}$ which correspond to an eigenvalue (winding number argument). Clearly gives lower bound.
- Upper and lower bound agree!!!!


## Theorem

If $Q_{\text {top }}=0$ and $u_{0} \in[0, \pi)$, let $N$ be the largest non-negative integer such that

$$
\int_{-\infty}^{\infty} \sin \left(\frac{u}{2}\right) d x>(2 N-1) \pi
$$

Then there exists exactly $N$ eigenvalues of (1), all of which are on the unit circle and are simple.

## Possible Extensions

- NUMERICS: Tentative numerical experiments indicate the result for $Q_{\text {top }}= \pm 1$ is tight: monotone kinks of higher topological charge and non-monotone kinks of topological charge $\pm 1$ frequently have eigenvalues off $S^{1}$.


## Possible Extensions

- NUMERICS: Tentative numerical experiments indicate the result for $Q_{\text {top }}= \pm 1$ is tight: monotone kinks of higher topological charge and non-monotone kinks of topological charge $\pm 1$ frequently have eigenvalues off $S^{1}$.
- Similar experiments on breather like potentials suggests this result may be improved to the case $0<u_{0}<2 \pi$. This is further supported by the fact that the winding number in the symplectic group is monotone until $u_{0}=2 \pi$.


## Possible Extensions

- NUMERICS: Tentative numerical experiments indicate the result for $Q_{\text {top }}= \pm 1$ is tight: monotone kinks of higher topological charge and non-monotone kinks of topological charge $\pm 1$ frequently have eigenvalues off $S^{1}$.
- Similar experiments on breather like potentials suggests this result may be improved to the case $0<u_{0}<2 \pi$. This is further supported by the fact that the winding number in the symplectic group is monotone until $u_{0}=2 \pi$.
- Thank You!

